

The Role of Jets in Black-hole X-ray Binaries

Nick Kylafis
University of Crete

with
I. Papadakis, P. Reig, D. Giannios, G. Pooley

Agios Nikolaos, 5 May 2010

Introduction

- In my opinion, the jets from compact X-ray sources have not been given proper attention.
- Most people treat jets simply as “fireworks”, which do nothing else than emit radio waves.

I hope to convince you that

- The jet is a central player in the observed phenomena and not simply an embellishment.

The jet model

- In a series of four papers

Reig, Kylafis, Giannios 2003, A&A

Giannios, Kylafis, Psaltis 2004, A&A

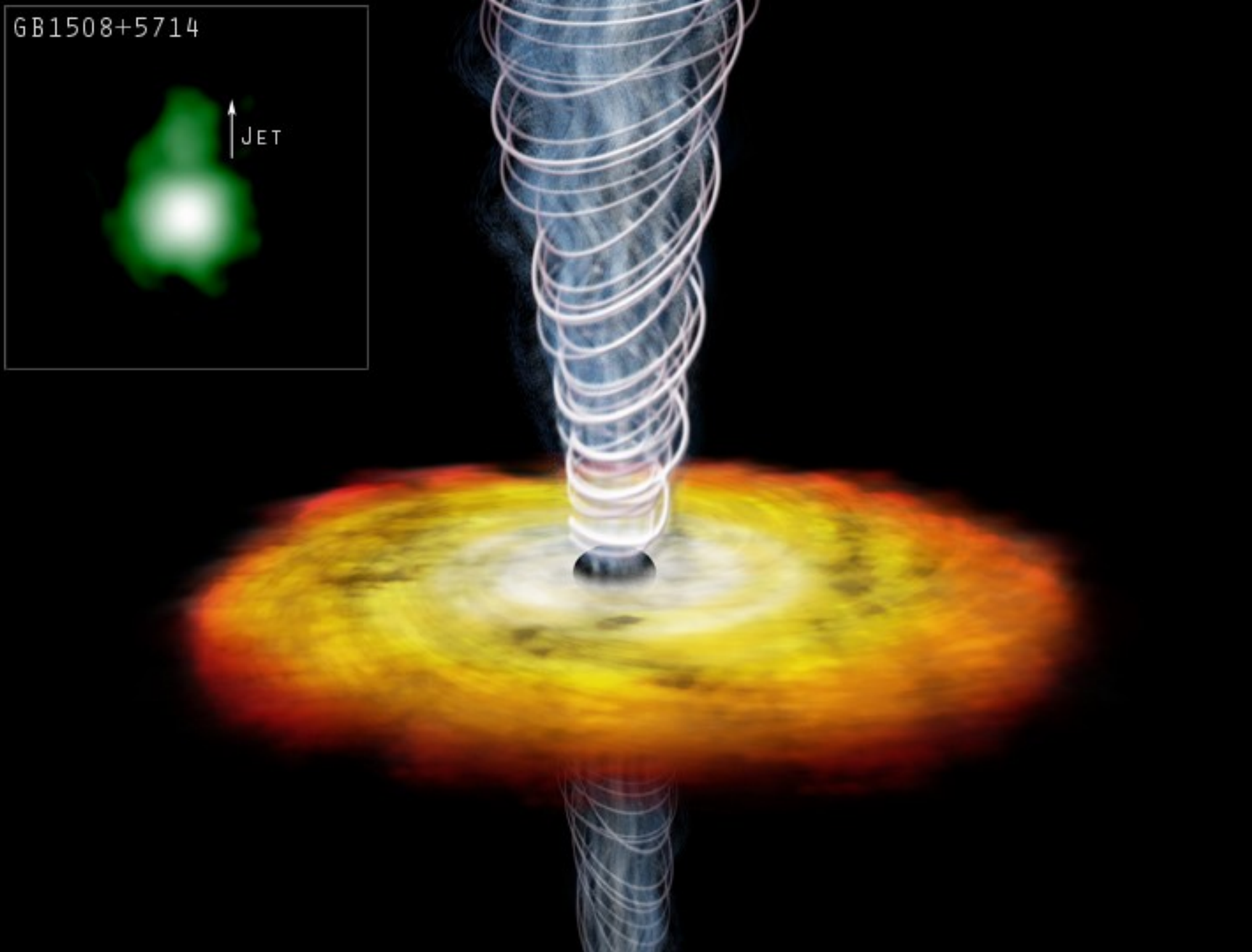
Giannios 2005, A&A

Kylafis, Papadakis, Reig, Giannios, Pooley, 2008

we proposed a jet model that explains a number of observational facts, when the black-hole X-ray sources are in the hard state.

GB1508+5714

JET



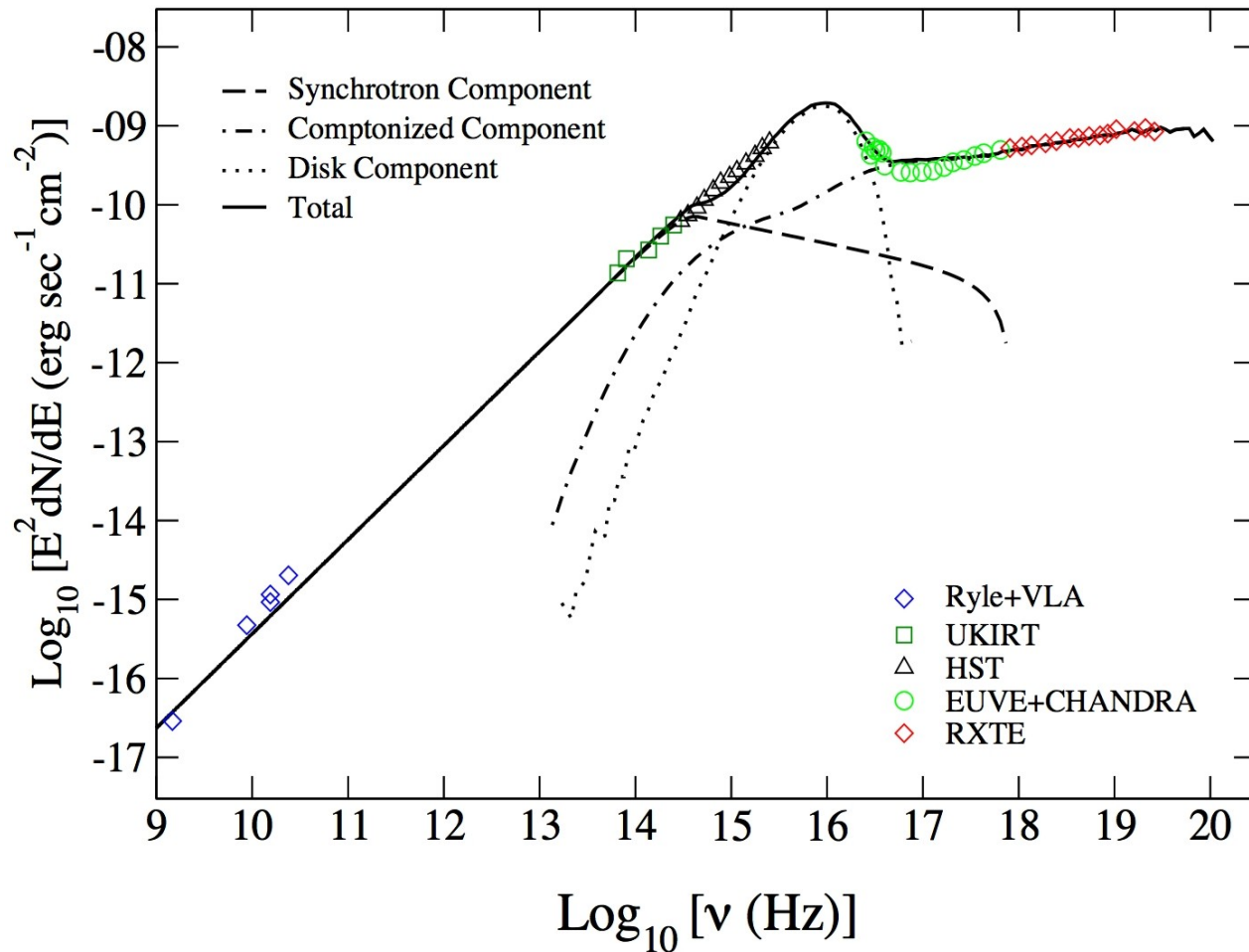
Ingredients of the model

- The jet is semi-relativistic ($v \sim 0.8 c$).
- The density in the jet falls off inversely proportional to distance from the black hole. Such flows are allowed theoretically (Vlahakis & Koenigl).
- In the rest frame of the flow, there is a power-law distribution of electron γ 's (standard assumption for radio jets).
- Soft photons from the accretion disk get up-scattered in the jet and a power-law spectrum is produced in hard X-rays (photon number index Γ).

Observational facts

- ENERGY SPECTRUM
- Up to now, only for one source (XTE J 1118+480) we have simultaneous observations from radio to hard X-rays.

Giannios 2005, A&A



- Model and observations for XTE J 1118+480

Life is not so easy however ☹️

- ❑ Impressive as the model fit may be, it DOES NOT constrain the model!
- ❑ Equally good fits to the data are produced by other models:
 - Markoff, Falcke, Fender 2001, A&A
 - Vadawale, Rao, Chakrabarti 2001, A&A
 - Corbel & Fender 2002, ApJ
 - Markoff et al. 2003, A&A

Let's see why.

Comptonization

- ❑ One can get a power-law X-ray spectrum with either thermal or non-thermal electrons.
- ❑ Also, the shape of the Comptonizing region can be spherical, cylindrical, planar, whatever!
- ❑ Thus, the energy spectrum alone CANNOT constrain the model.

Time lag between hard and soft X-rays

- It has been observed (Nowak et al. 1999; Ford et al. 1999) that the hard X-rays (say 8 -14 keV) lag the soft X-rays (say 2 - 4 keV).
- This is expected for Comptonization models.
- However, the time-lag is a function of Fourier frequency!!!
- This is contrary to intuition for simple Comptonization models.

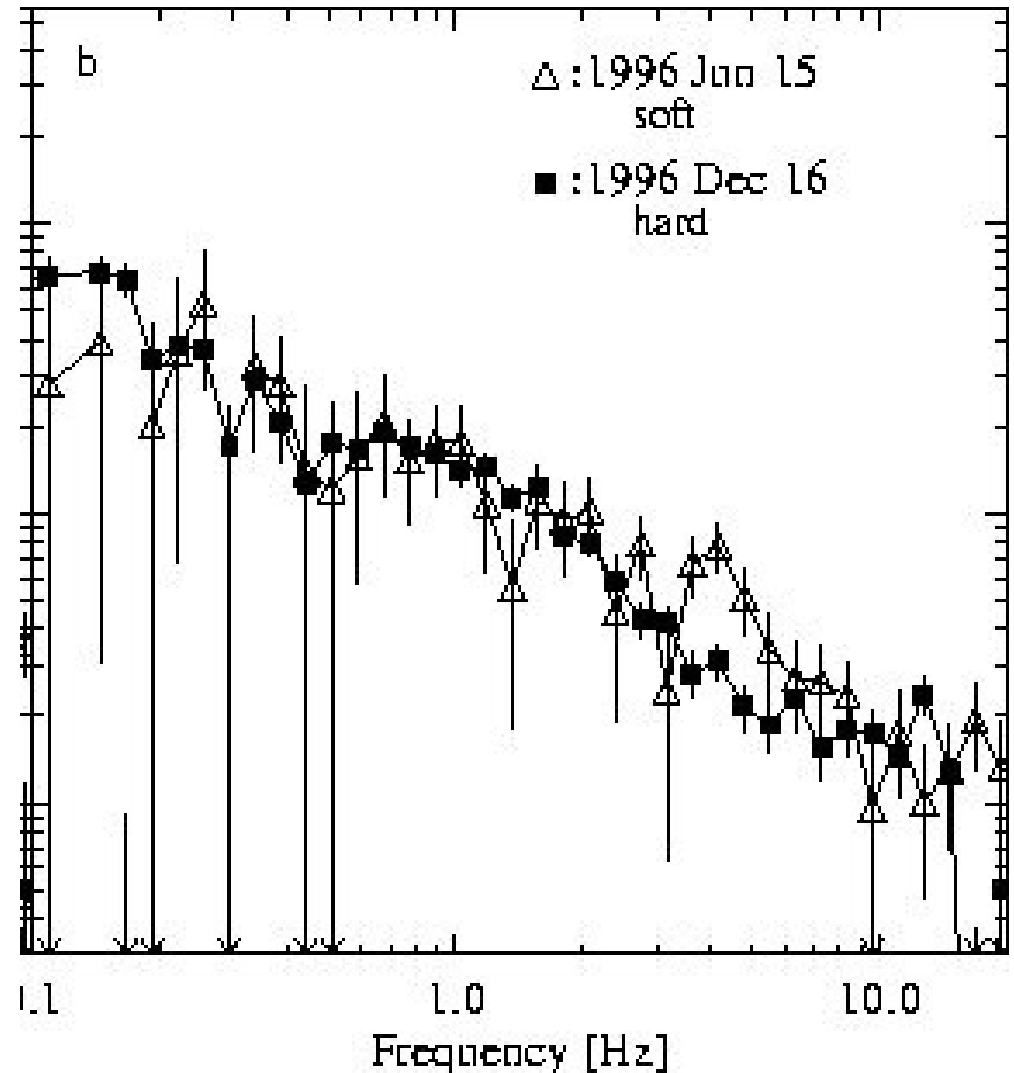
Time lag between hard and soft X-rays

- For Cyg X-1,

$$t_{lag} \propto \nu^{-\beta}, \beta \approx 0.7$$

Time lag vs Fourier frequency

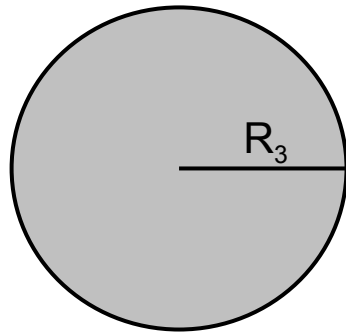
Time lag



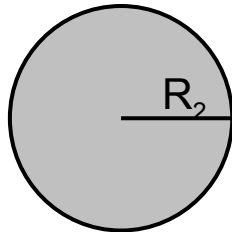
Compton scattering acts like a filter

- It cuts off the high frequencies.
- If (period of variability) $<$ (time lag), the variability is washed out.
- Therefore, frequencies $> 1/(\text{time lag})$ are not observed.

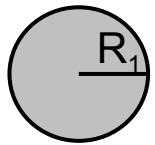
Schematic picture of our jet



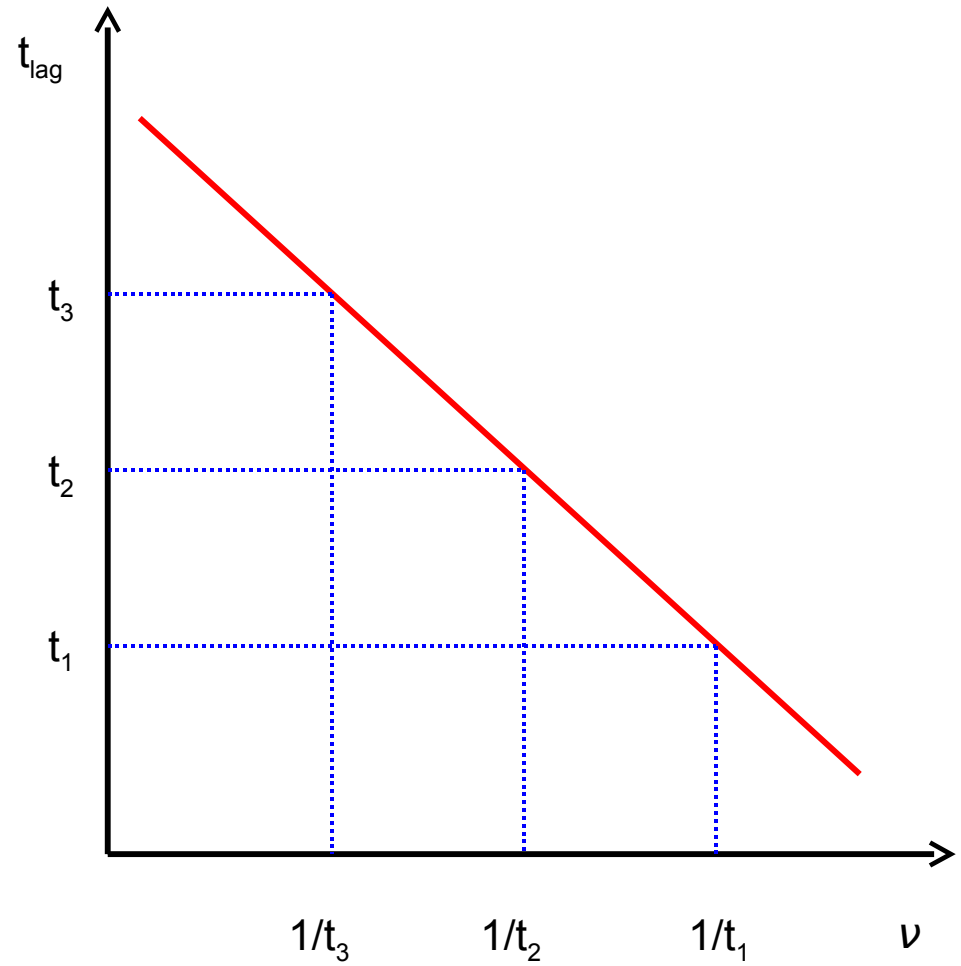
t_3



t_2



t_1



Other models

- Constraining as it is, the time lag vs Fourier frequency relation has been explained by other models also:

Poutanen & Fabian 1999, MNRAS

Kotov, Churazov, Gilfanov 2001, MNRAS

Koerding & Falcke 2004, A&A

Therefore, even more constraints are needed.

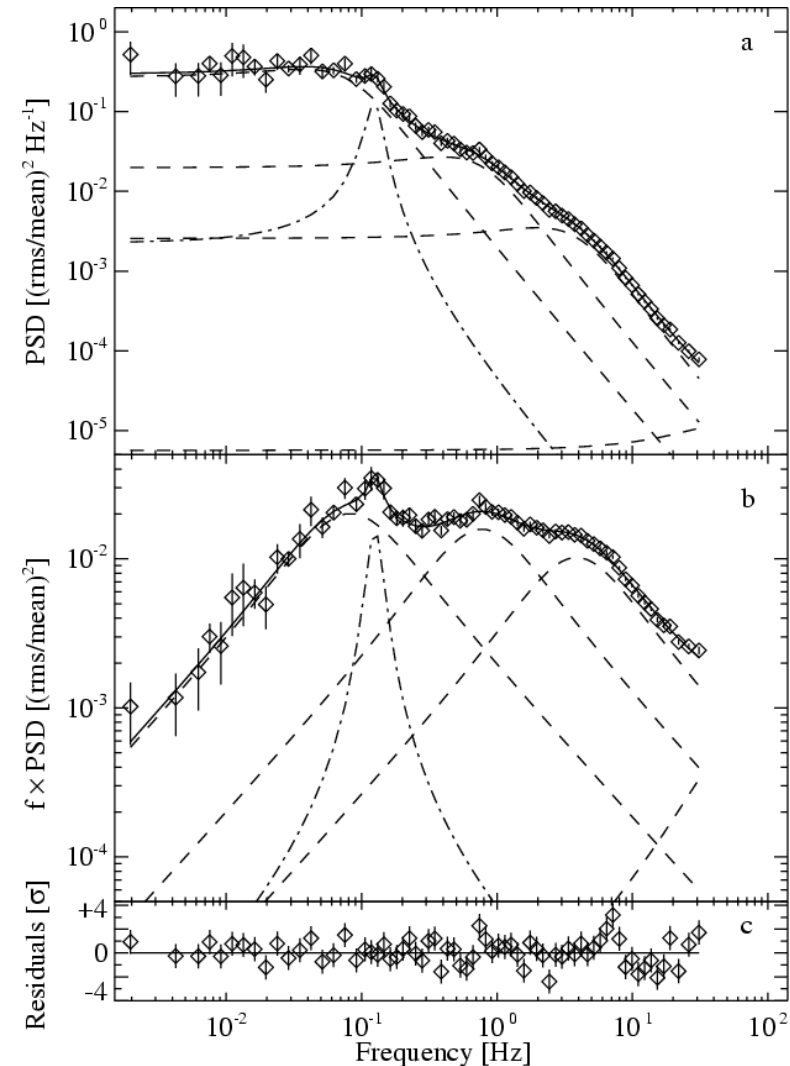
More observational constraints

- The **long-term variability** of Cyg X-1 has been studied by Pottschmidt et al. (2003), A&A (see also Shaposhnikov & Titarchuk 2006; 2007).
- When the source was in the hard state, the study revealed a number of very stringent constraints.
- These are:

Pottschmidt et al. (2003), A&A

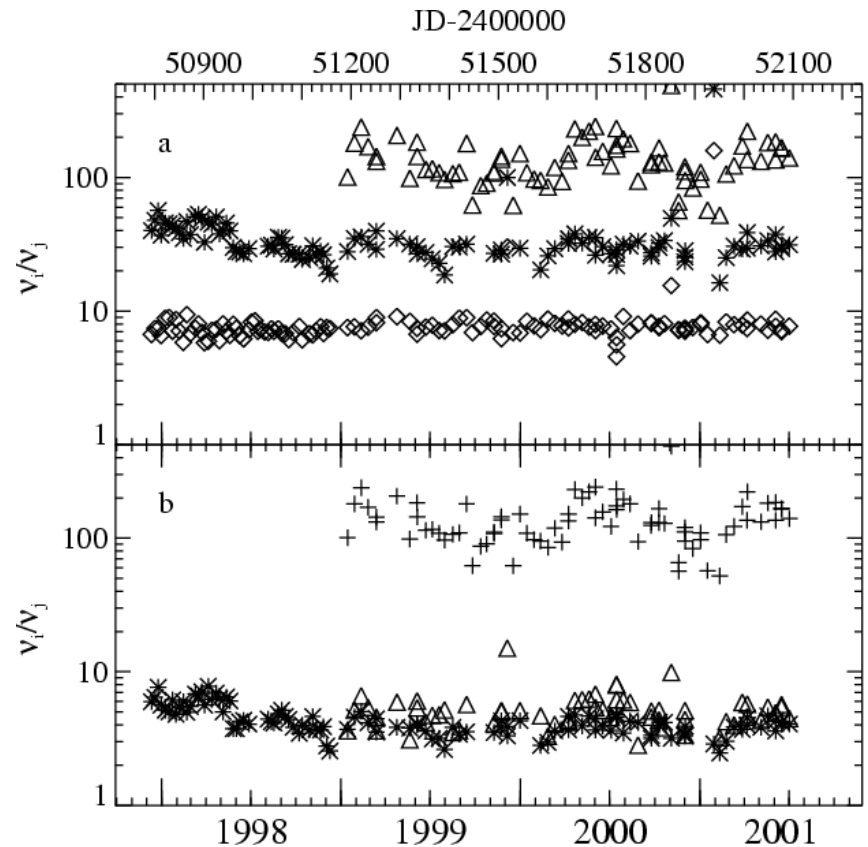
- The power spectrum of Cyg X-1 was fitted with four broad Lorentzian profiles that have peak frequencies

$$\nu_1, \nu_2, \nu_3, \nu_4$$



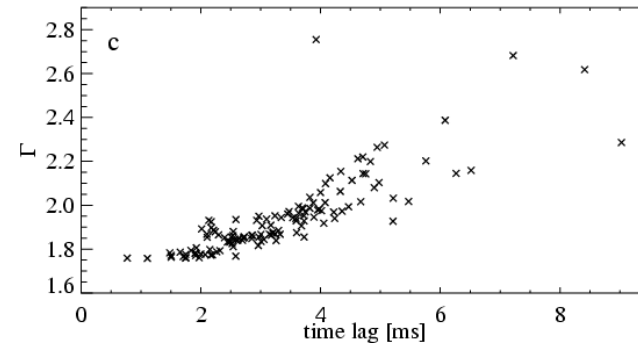
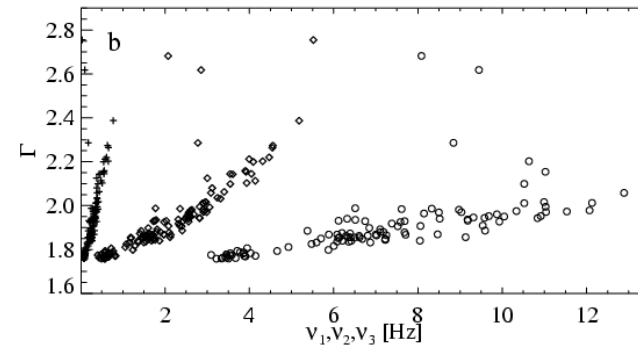
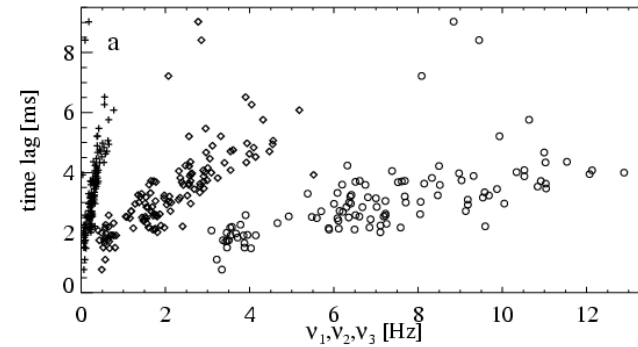
Pottschmidt et al. (2003), A&A

- The ratios of the peak frequencies are CONSTANT!!!



Pottschmidt et al. (2003), A&A

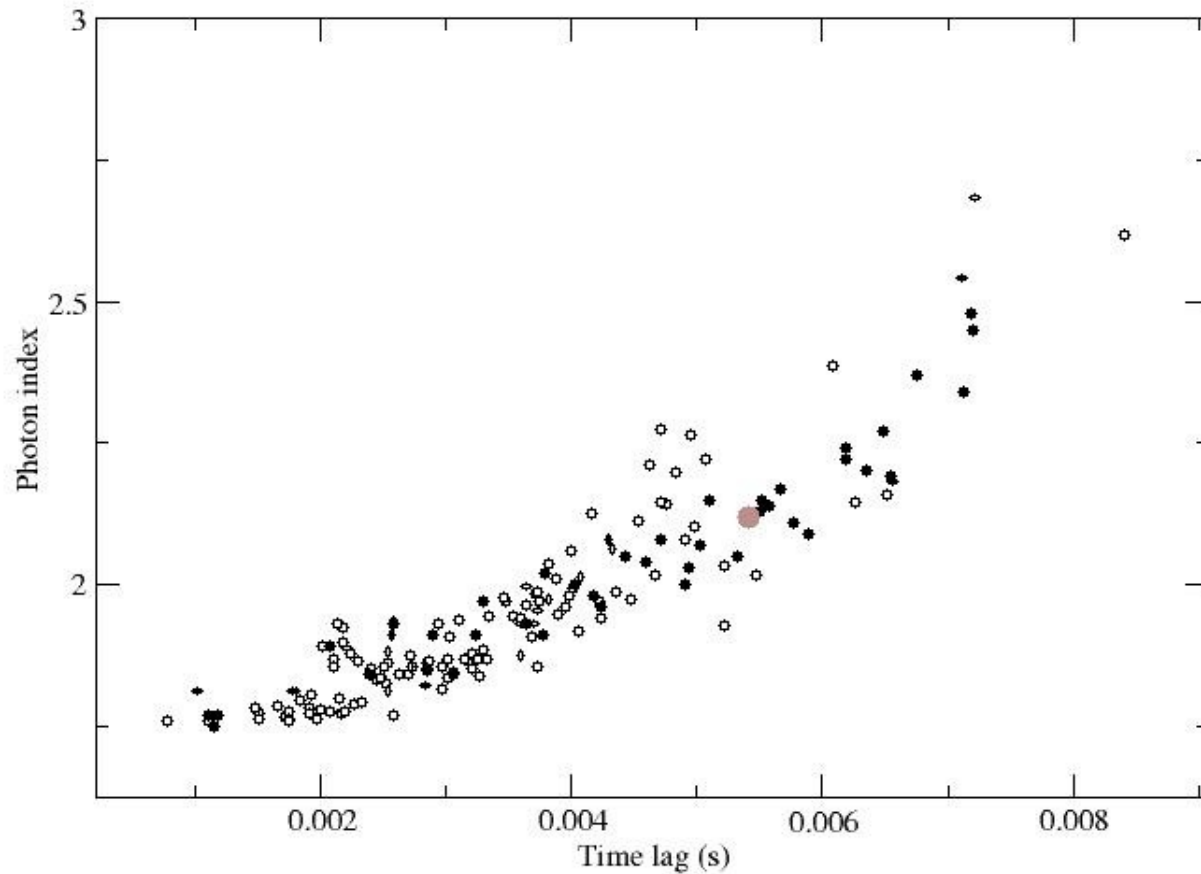
- Very stringent correlations:



Our jet model

- Question: Can we explain these correlations by simply varying the parameters of our model around their typical values?
- The answer is YES! (else I would not pose the question).
- We varied only two parameters: The density (or equivalently the optical depth) and the radius of the base of the jet.
- We were thus able to reproduce the Γ - $\langle \text{timelag} \rangle$ correlation.

Gamma vs. $\langle \text{timelag} \rangle$



Identification of the Lorentzian peak frequencies.

- Using just the density and the radius at the base of the jet, can we think of a combination that has the dimensions of frequency (inverse timescale)?

Identification of the Lorentzian peak frequencies.

- The only inverse timescale, that we can think of, is

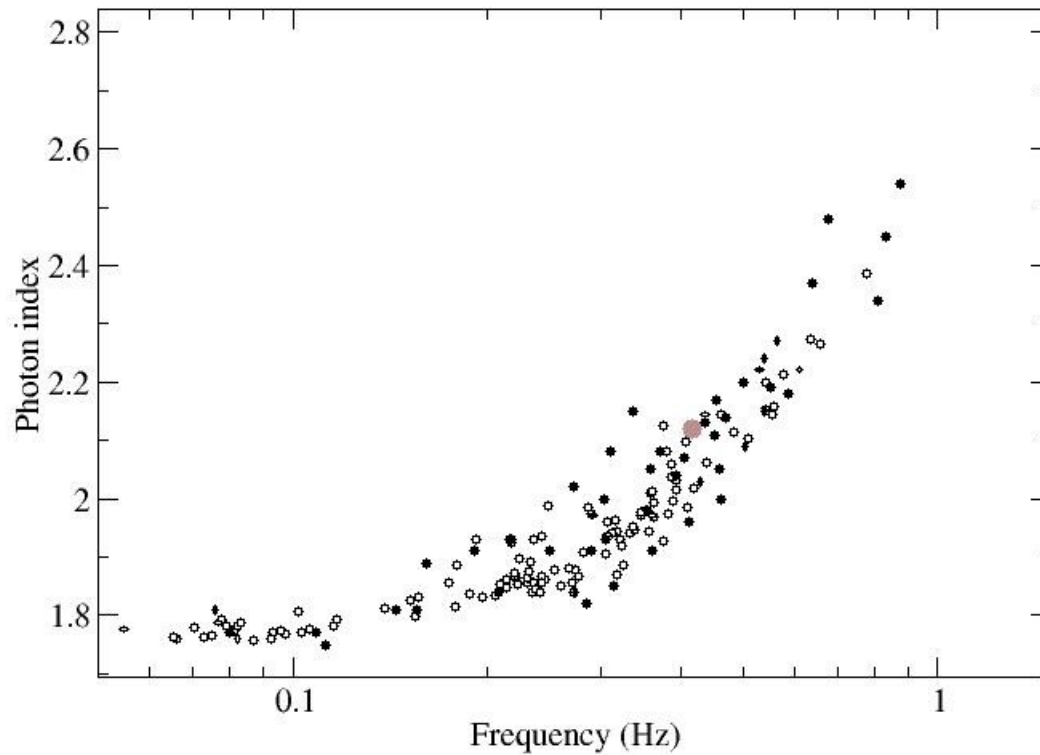
$$\frac{\dot{M}_{outflow}}{M_{available\ for\ outflow}}$$

Identification of the Lorentzian peak frequencies

$$\frac{\dot{M}_o}{M_o} = \frac{\pi R_o^2 n_o u_{flow}}{M_o} = CR_o^2 n_o \propto v_1$$

Gamma vs. peak frequency

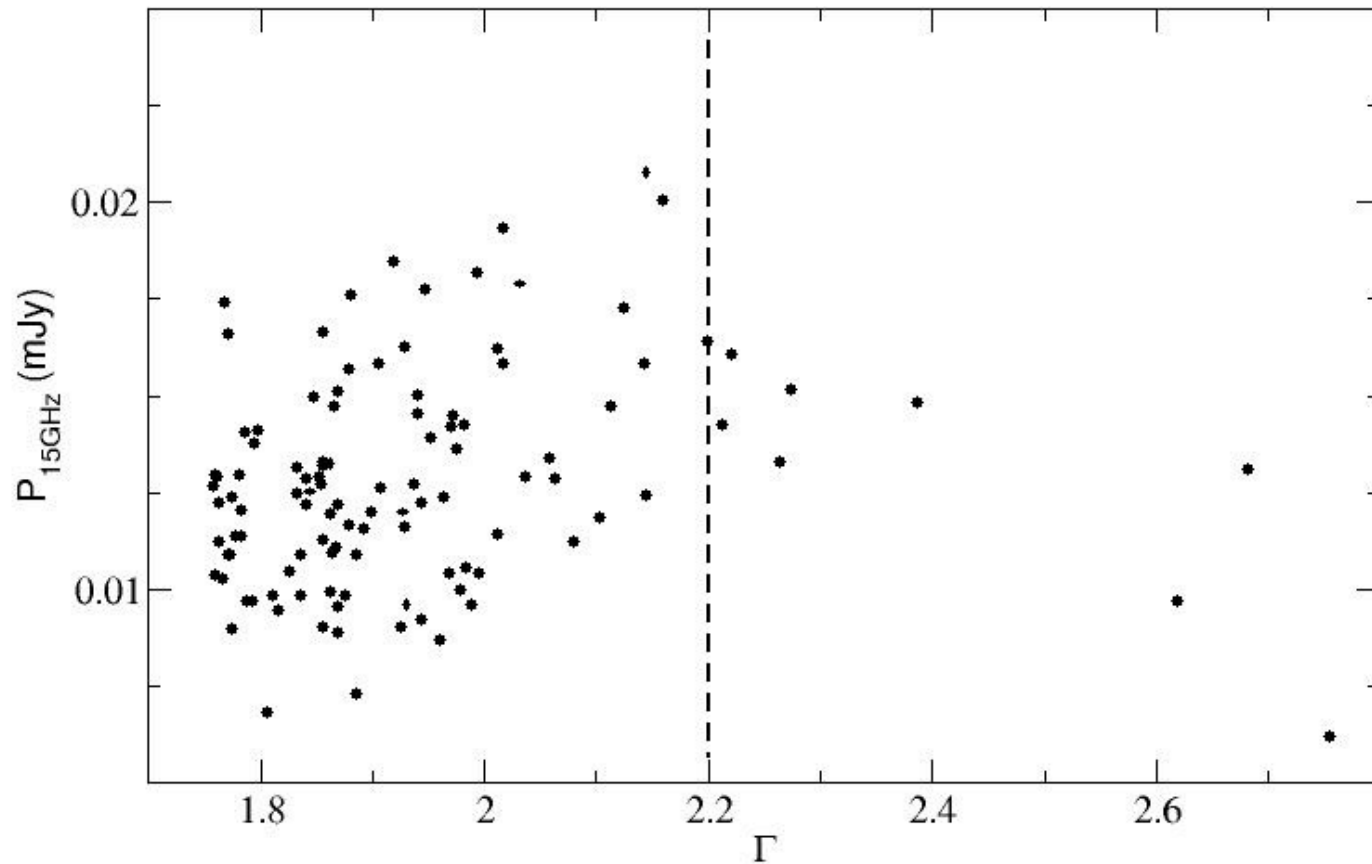
1



Additional constraint

- Our jet model predicts a positive correlation between radio flux and Γ .
- Such a correlation has not been seen or proposed before.
- In Cyg X-1 we have found this correlation. It is not very tight, but it is certain (Kendall's tau=0.21, i.e., probability $< 0.2\%$ that there is no correlation).

Radio flux vs Gamma



Additional constraints

- Our jet model **must** predict/confirm the reflection component and its correlations with other observables.
- It must also explain the 0.1 s lag of the IR with respect to X-rays in GX 339-4 (Casella et al. 2010).
- Our simplifying, but unphysical assumption, that the jet velocity is constant throughout the jet must be modified. **The acceleration of the jet must be taken into account.**

Additional constraints

- Simple estimates suggest that we can explain both the time lag and the reflection component.
- We must make sure though that the introduction of an acceleration region at the base of the jet does not affect all our previous results.
- We are currently working on it.
- Thanks